

Physics of varying spacetime dimensions

Moorea, PACIFIC 2024

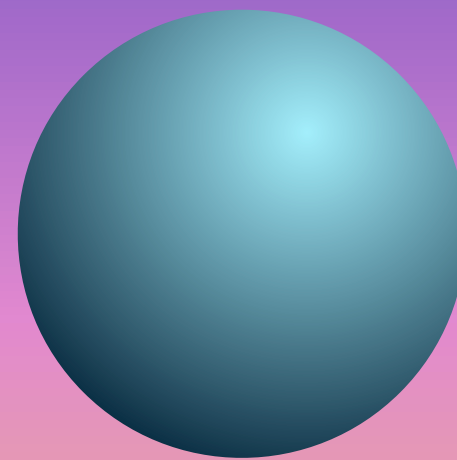
Tom Melia, Kavli IPMU

Cao, Lu, TM SciPost Phys. Core 7, 055 (2024)

Q:

What's the
2D version of
a sphere?

2D



sphere

3D

Circle



2D

sphere



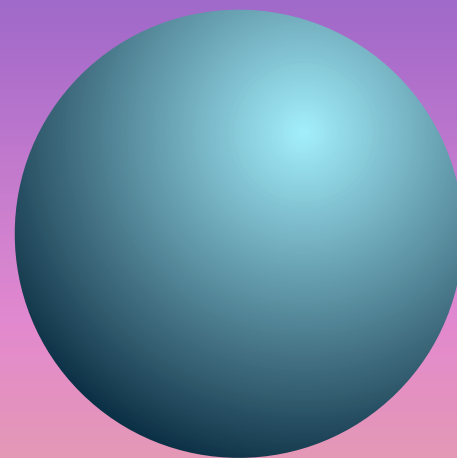
3D

Circle



2D

sphere



3D

Q:

What's the
4D version of
a sphere?

4D

Circle



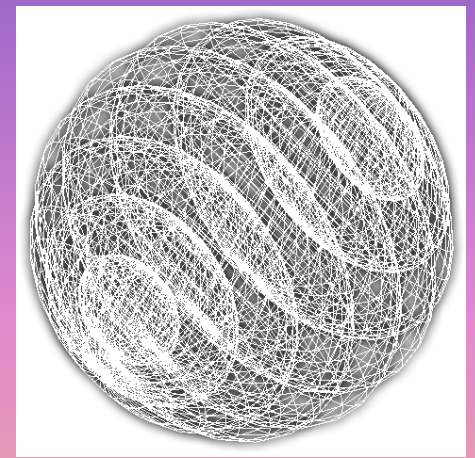
2D

sphere



3D

Four dimensional
sphere



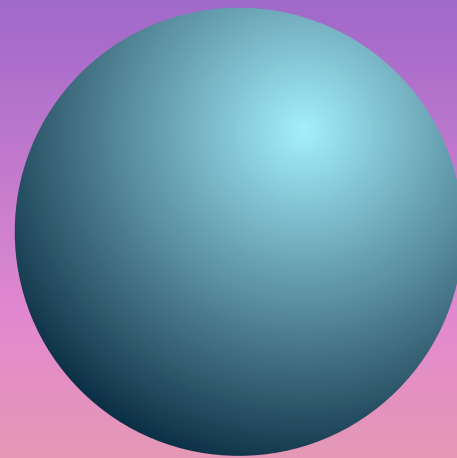
4D

Circle



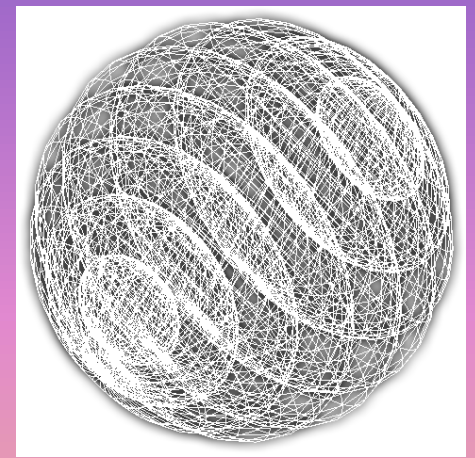
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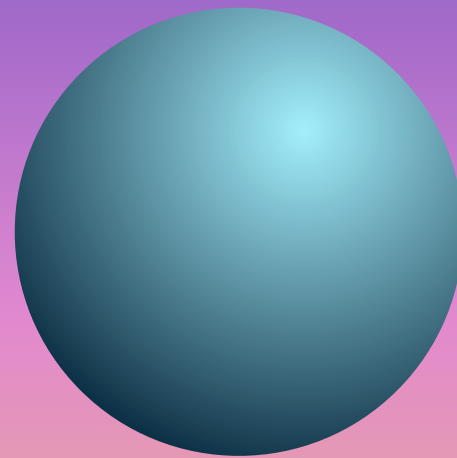
What's the
2.7434 D
version of a
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Circle



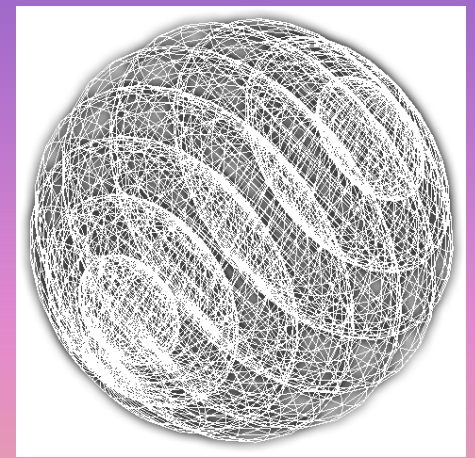
2D

sphere



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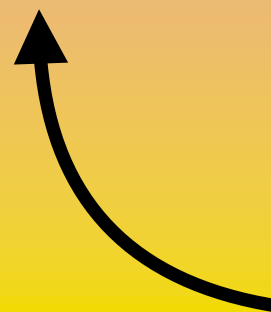
4D

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A:

A 2.7434
dimensional
sphere



D

Q:

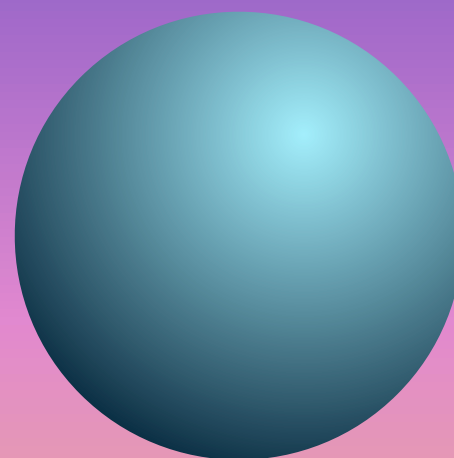
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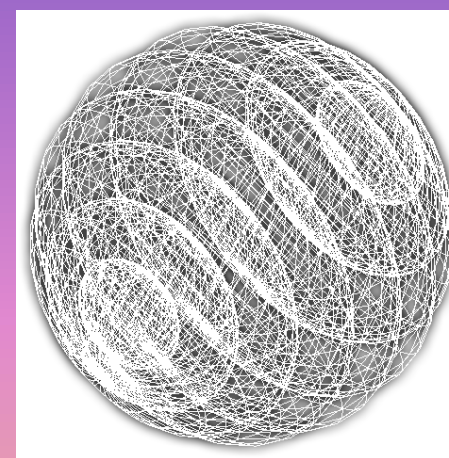
Circle



sphere



Four dimensional
sphere



1D

2D

3D

4D

D

Q:

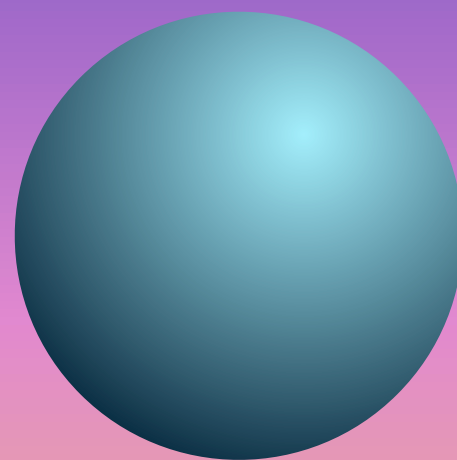
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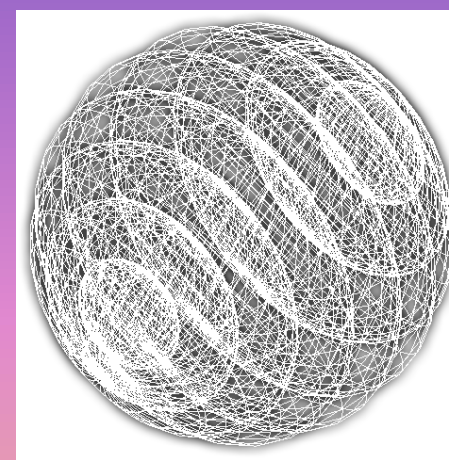
Circle



sphere



Four dimensional
sphere



1D

2D

3D

4D

D

A:

This talk

Q:

Why do we care?

Q: Why do we care?

A1: We here are all dimension ninjas

We think about higher/lower dimensions for all variety of phenomenological and theoretical reasons

We think about $4-\epsilon$ non-integer dimensions to do calculations (most of the time setting $\epsilon \rightarrow 0$ at the end, but not always, e.g. Wilson Fisher fixed point)

Do $1/N$ expansions

....

Q: Why do we care?

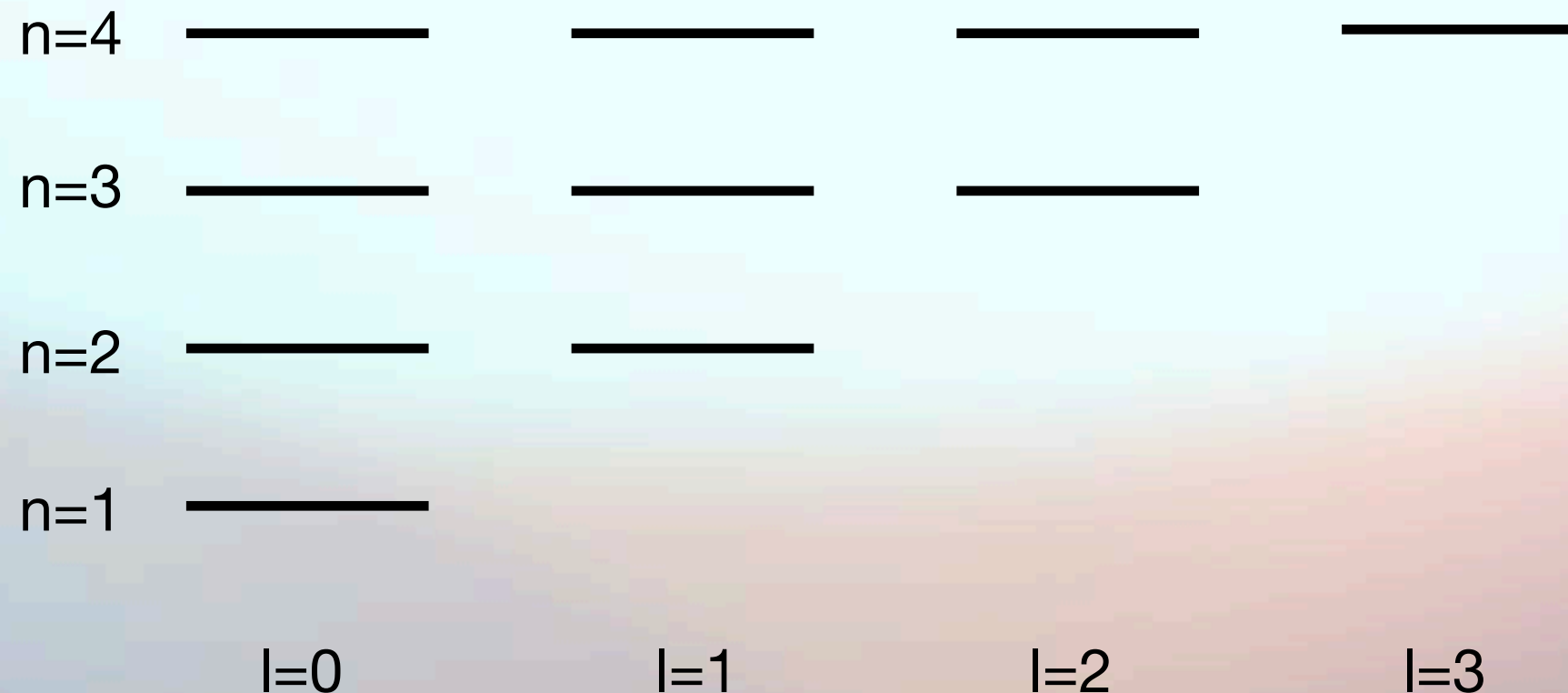
A2: Symmetry is the main tool we have to understand anything about QFT

Lets go back 100 years..

Degeneracy in QM

Quantum mechanics began when the spectrum of the Hydrogen atom was calculated, circa 1926 (Heisenberg, Pauli, Schrödinger)

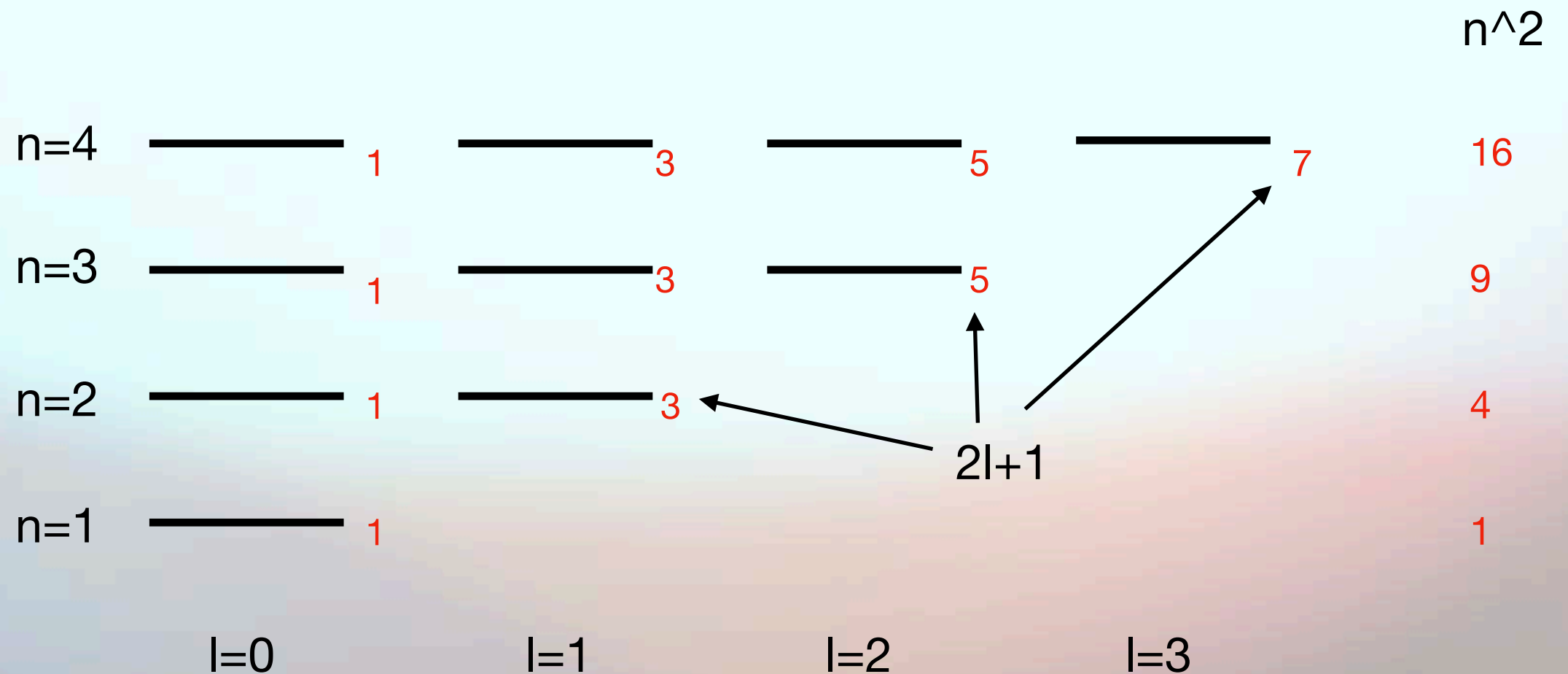
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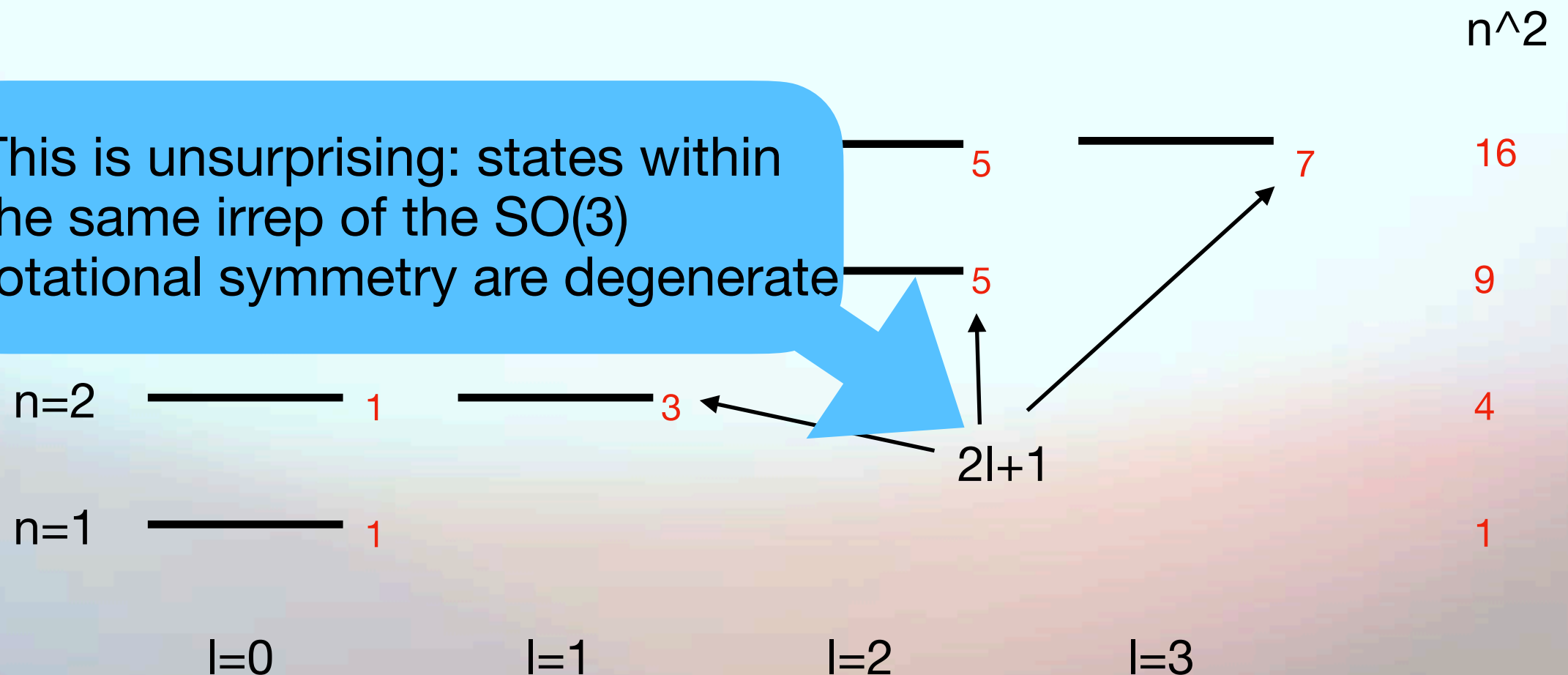


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$n=2$

$n=1$

$l=0$

The degeneracy between states of different $SO(3)$ irreps is down to a hidden $SO(4)$ symmetry

n^2

16

9

4

1

5

5

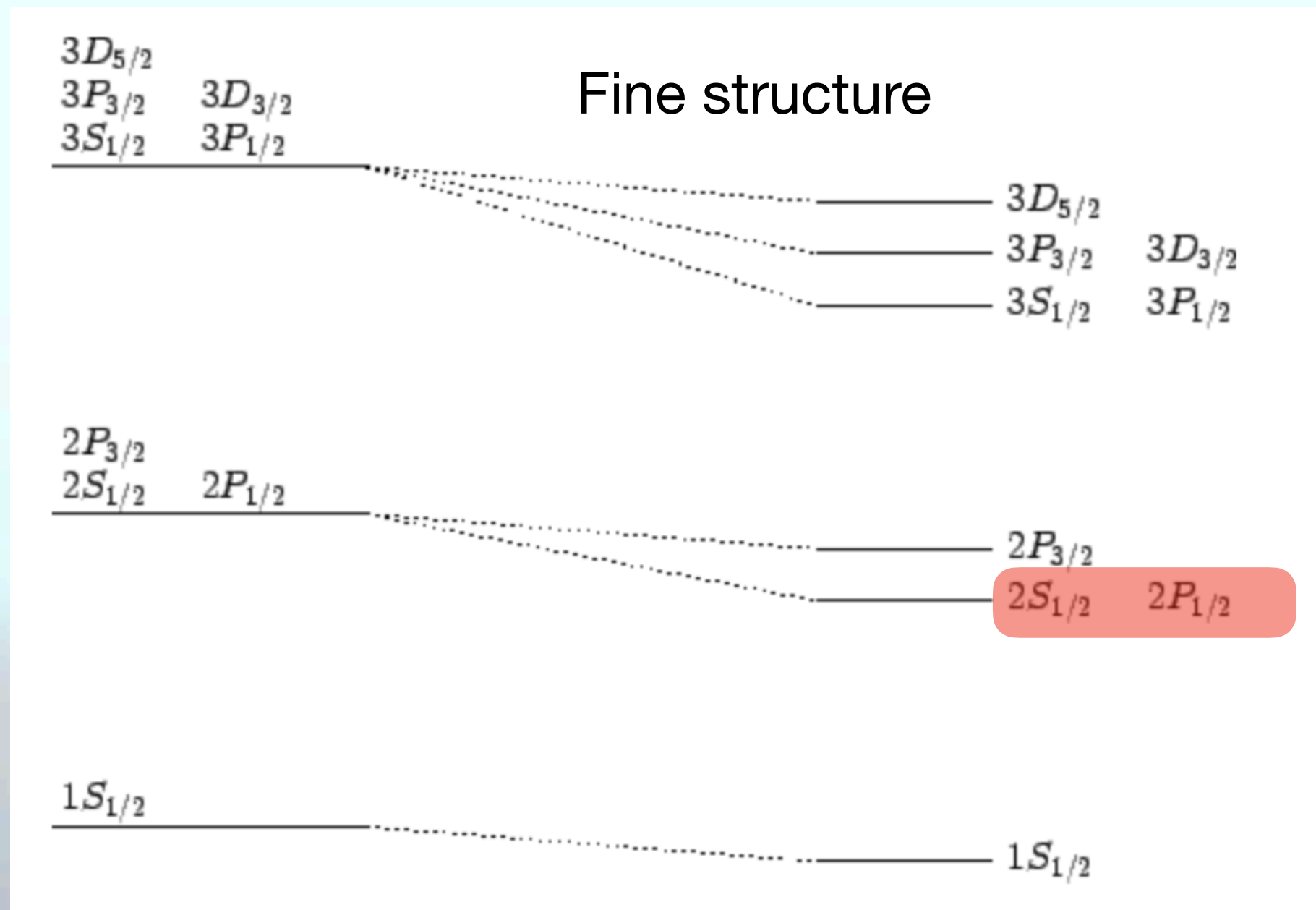
7

3

$2l+1$

Degeneracy in QM

The Lamb shift played a central role in the development of field theory



Level n
degeneracies
in first H
model

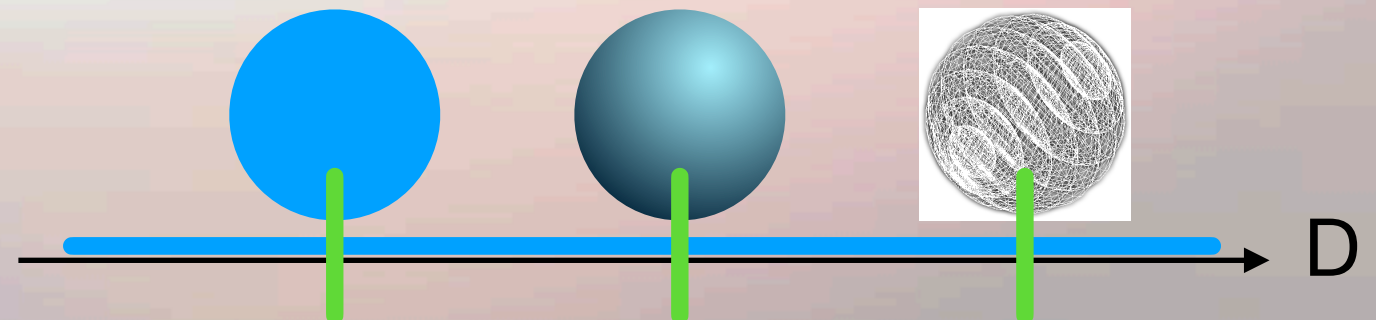
Degeneracies removed by small
perturbations

Level n
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Unitary QFT

Non-unitary QFT



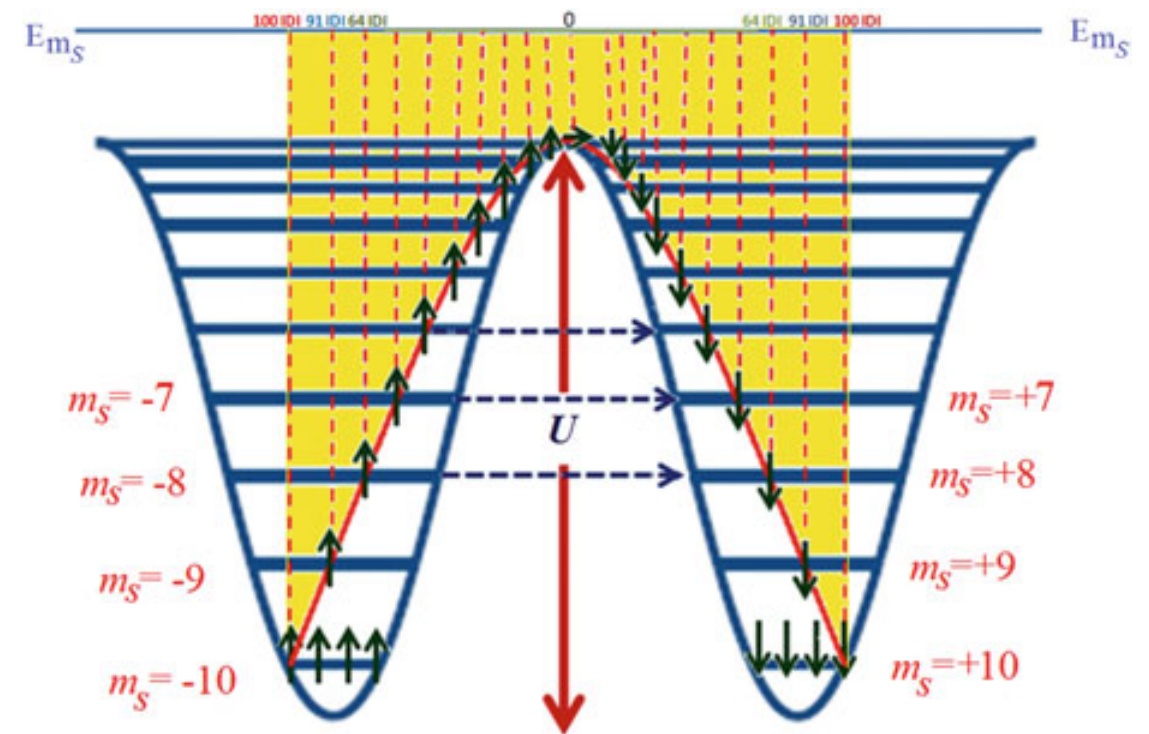
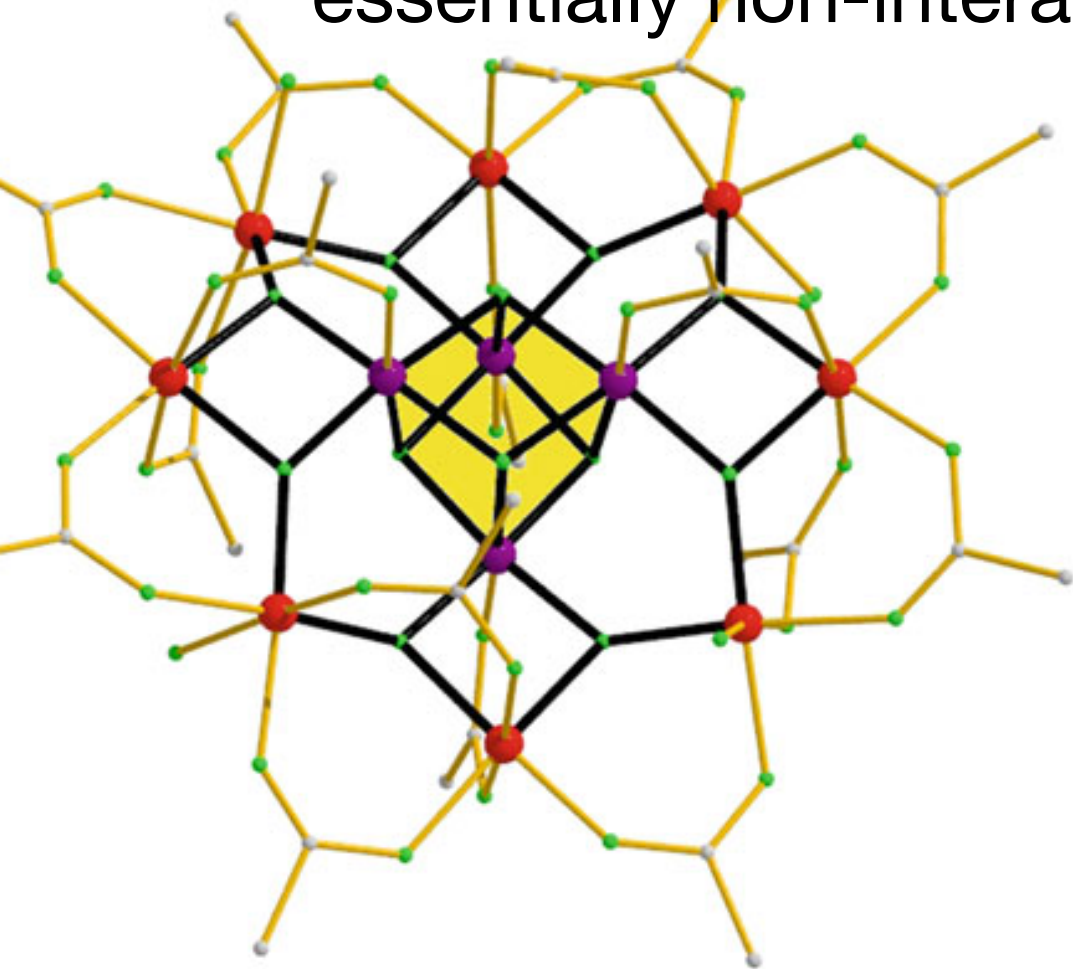
$O(N)$ for non-integer N : If the theory is a CFT, it is non-unitary

[Binder and Rychkov 1911.07895]

Degeneracy in QM

Macroscopic changes

Single molecule magnets are molecular crystals with essentially non-interacting spins at centre of each molecule

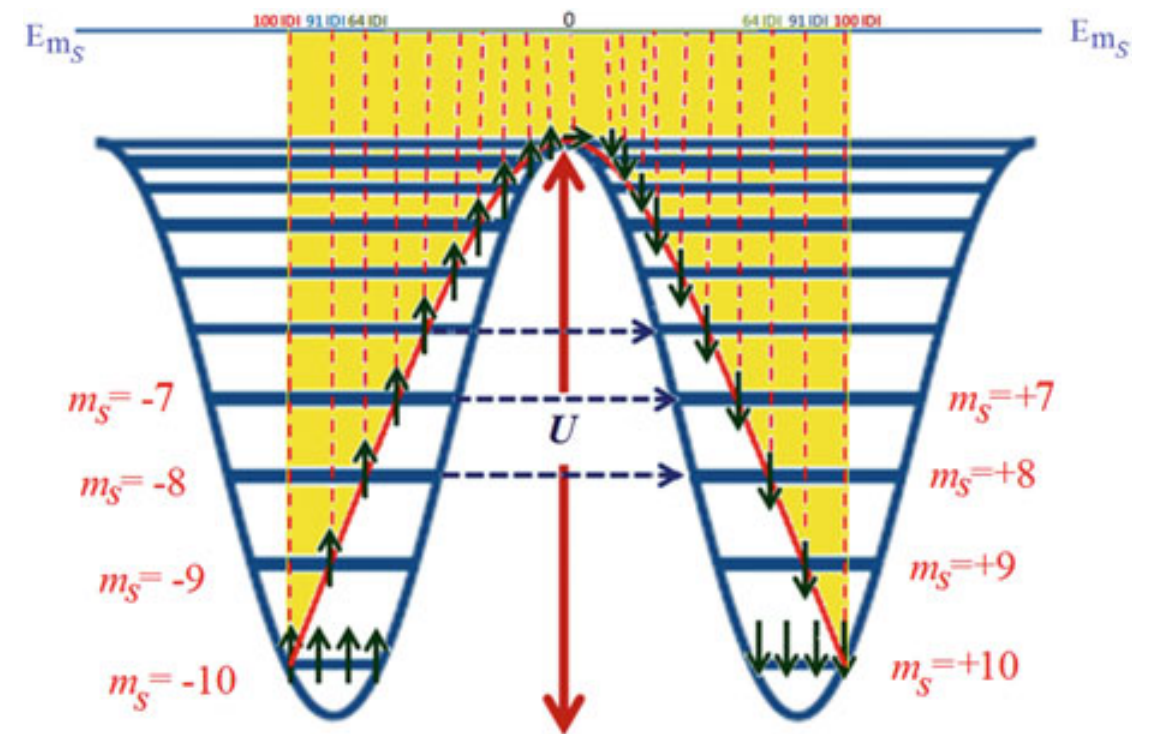
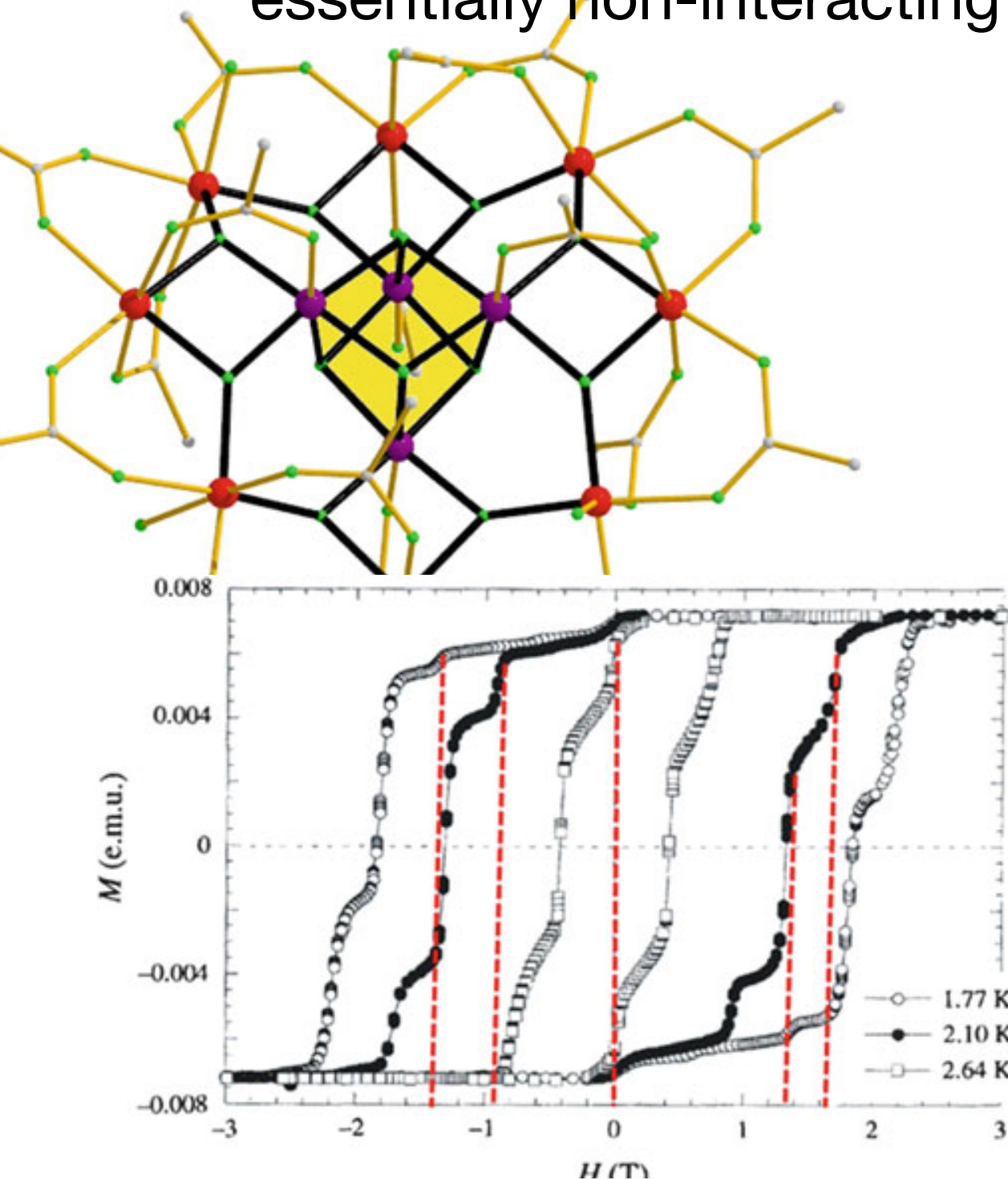


$$H = -U \sum_i (S_i^z)^2$$

Degeneracy in QM

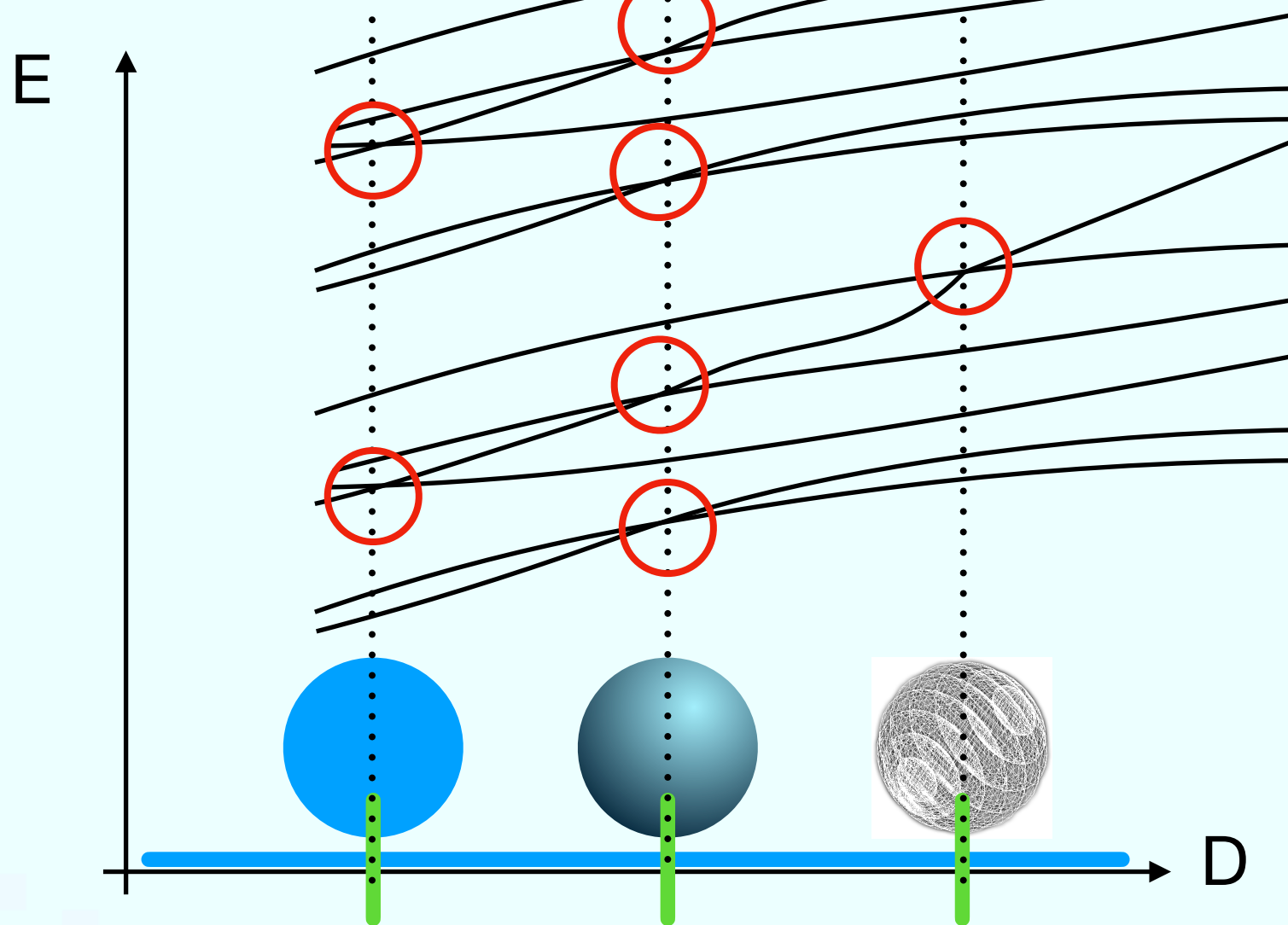
Macroscopic changes

Single molecule magnets are molecular crystals with essentially non-interacting spins at centre of each molecule



$$H = -U \sum_i (S_i^z)^2 - g S_i^z B$$

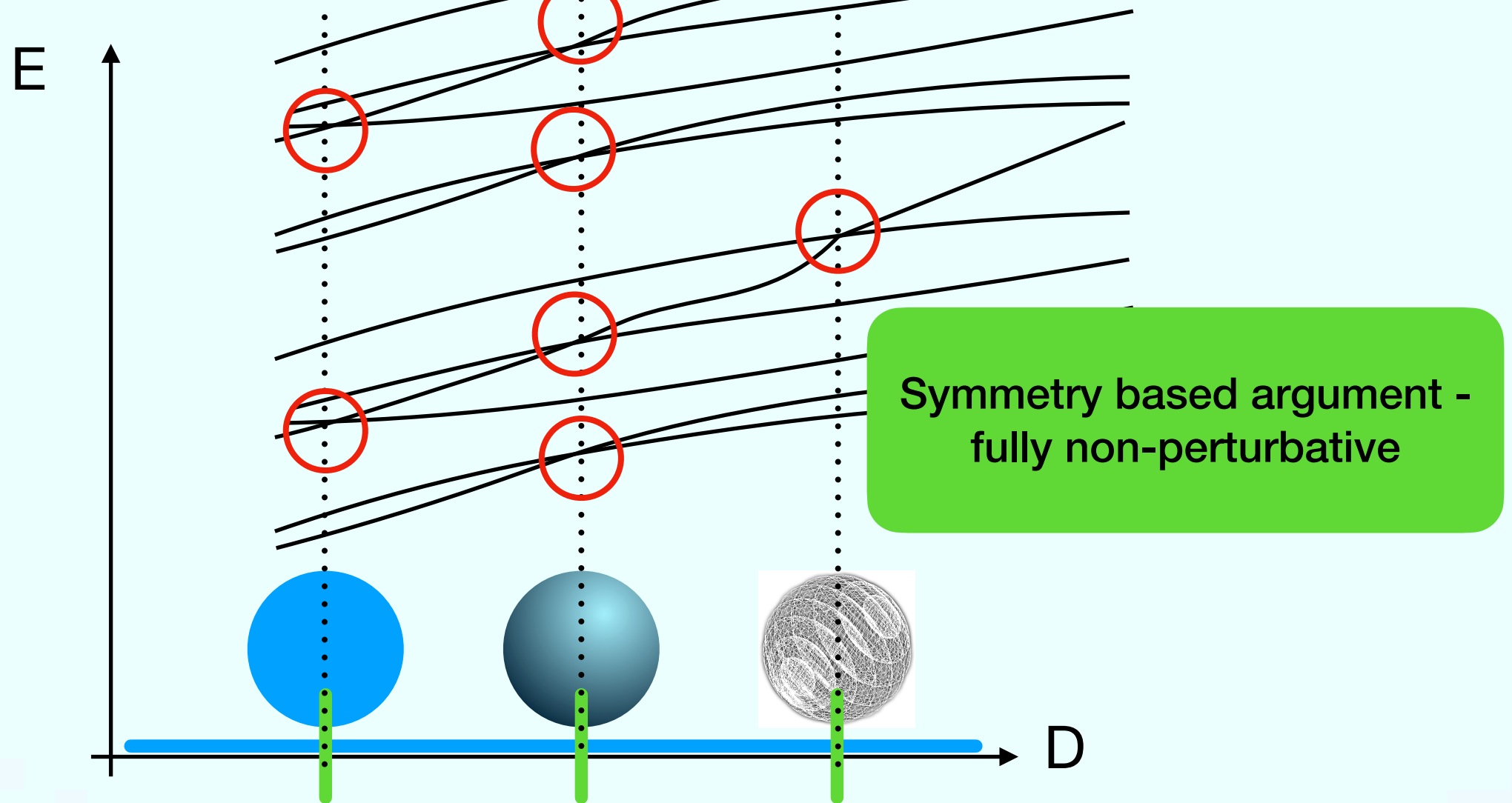
Tuning B field, degeneracies between the wells enhances quantum tunnelling, jumps in the hysteresis curve



Assume: spectrum continuity

Key result: the continued representation theory dictates some states drop out - “are evanescent” - *in pairs of equal energy*

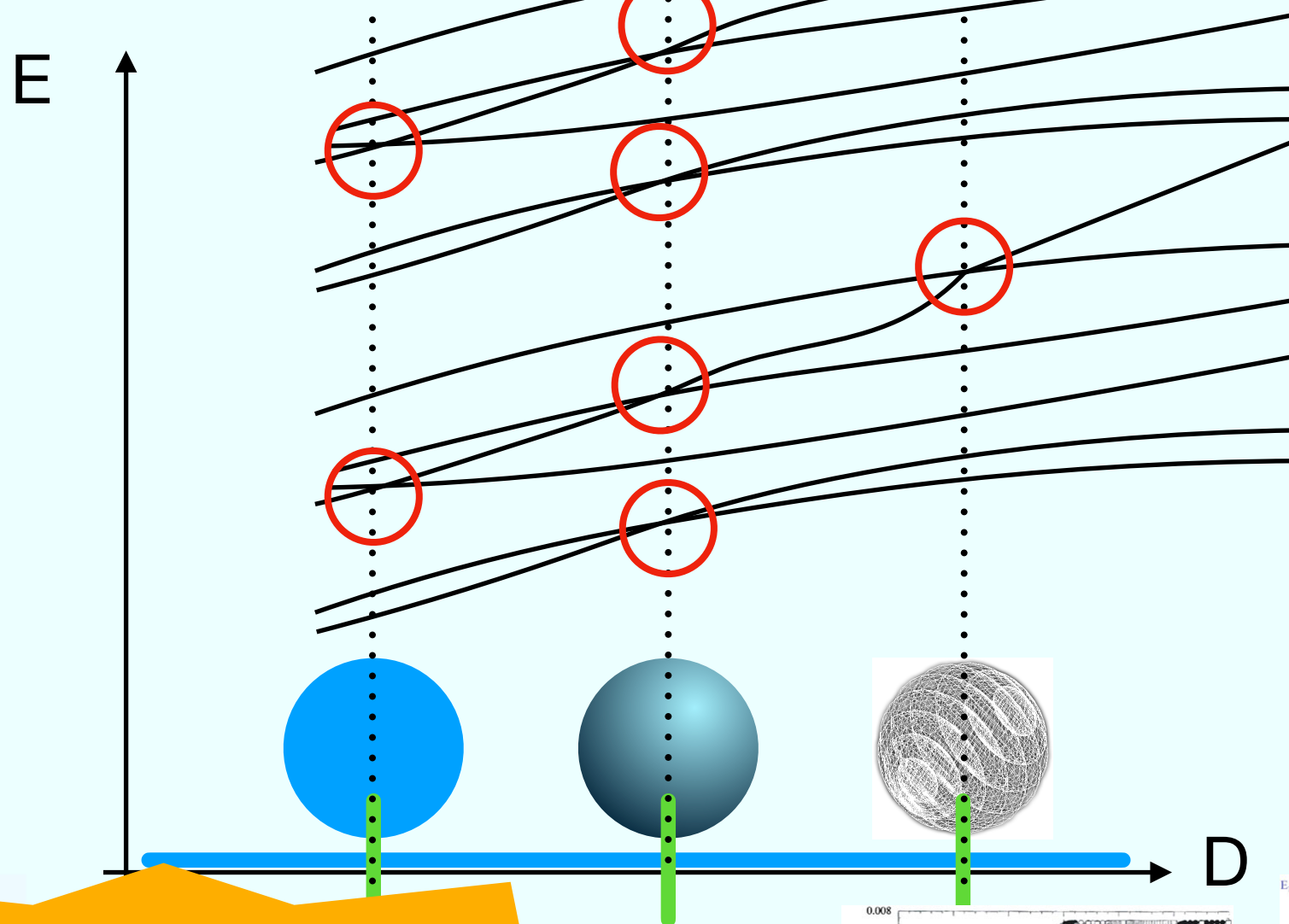
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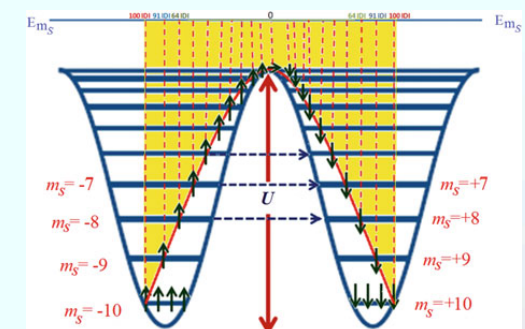
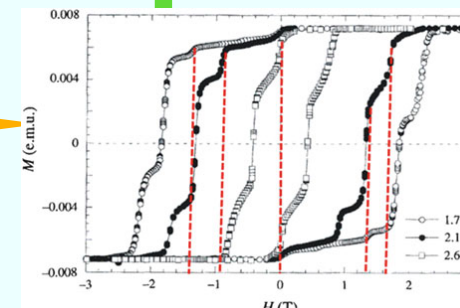
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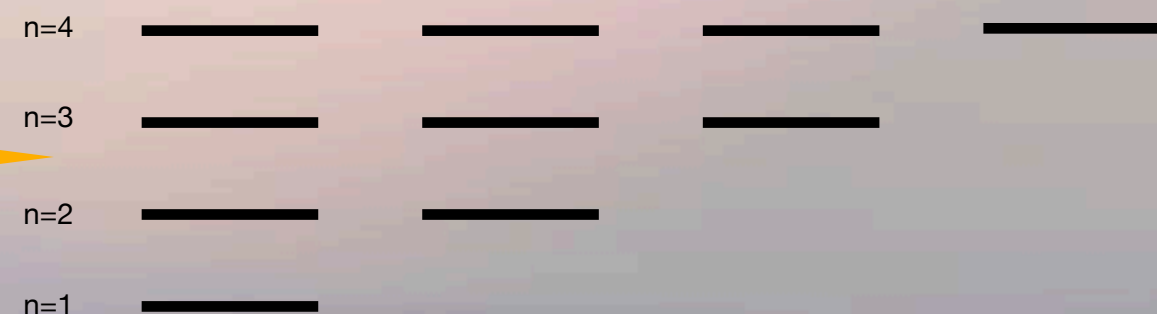
Degeneracy at isolated points of parameter N - when it is an integer

Accompanied by a major physical change: the theory becomes unitary

Occurs between different irreps of $O(N)$, a la Hydrogen spectrum



e.g. known existence of unitary islands with $N=1,2,3$ in $d=3$ via conformal bootstrap approach



Adding spins

We all learn $\frac{1}{2} \otimes \frac{1}{2} = 0 \oplus 1$

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Writing this in terms of SU(2) representations, we can use Young diagram

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...

SU(N>2): $\square \otimes \square = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} \oplus \square\square$

Adding spins

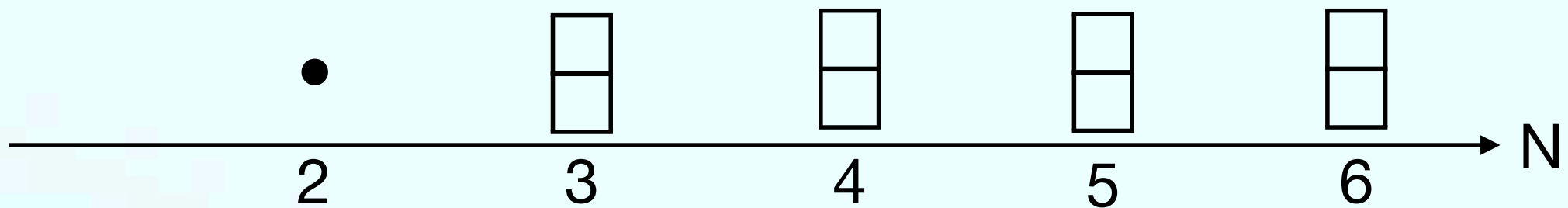
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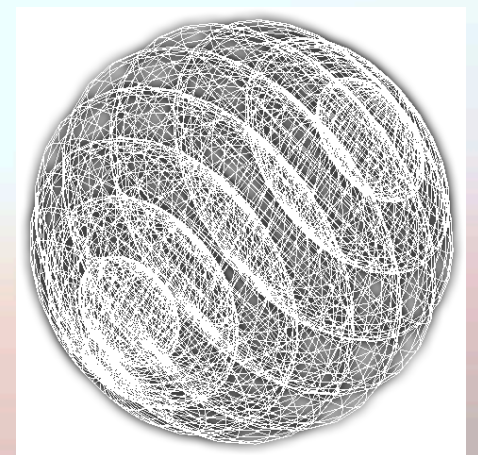
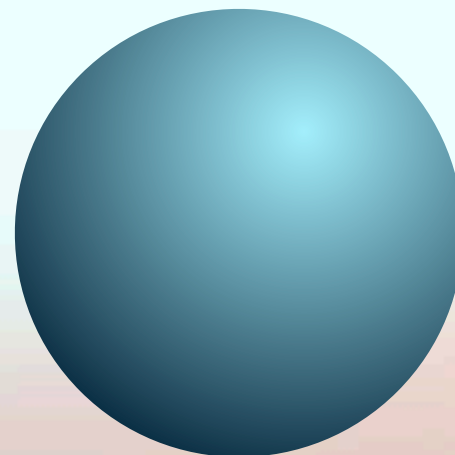
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c.f.

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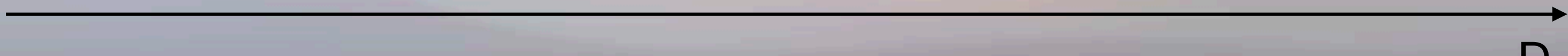


1D

2D

3D

4D

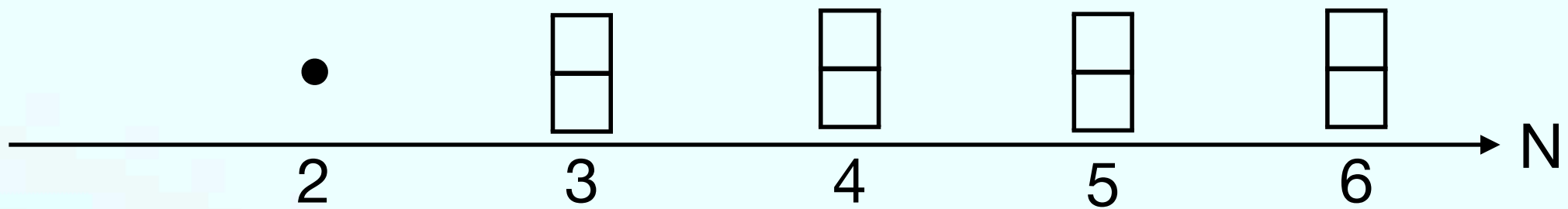


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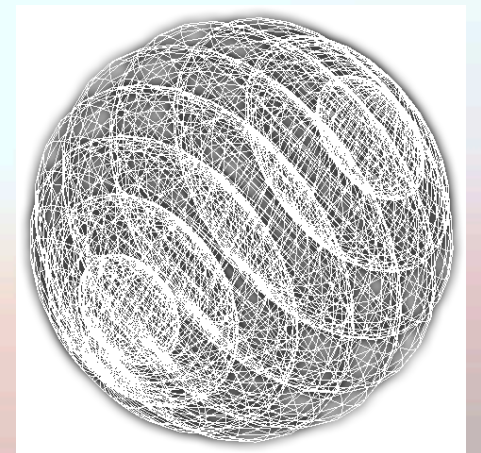
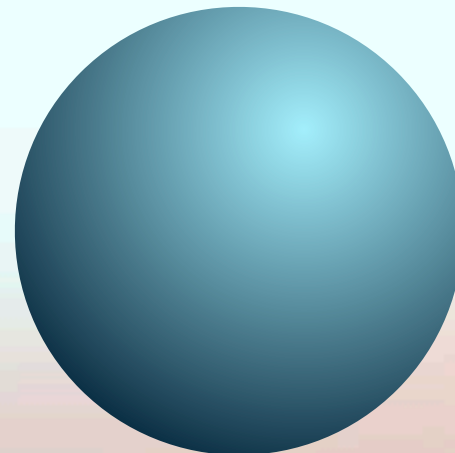
SU(N>2):

This is the right way to analytically continue.
Go to high enough N where we know what the
representation is.
Supplement with rules at special N



c.f.

?

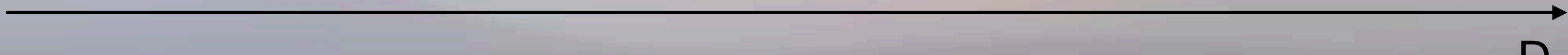


1D

2D

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4D



Analytic continuation for O(N)

Fixed integer N

$$R_1 \otimes R_2 = \sum_k c_k^{1,2}(N) R_k$$

Continued N

$$\bar{R}_1 \otimes \bar{R}_2 = \sum_k \bar{c}_k^{1,2} \bar{R}_k$$

With ‘large N saturated’ multiplicities and specialisation rules

$$\bar{c}_k^{1,2} = \lim_{N \rightarrow \infty} c_k^{1,2}(N)$$

$$\bar{R}_k \rightarrow ???$$

At ‘too low’ integer values of N

Analytic continuation for $O(N)$

Fixed integer N

These specialisation rules are known, but only relatively recently (in the context of the history of representation theory!)

K. Koike and I. Terada, Young-diagrammatic methods for the representation theory of the classical groups of type B_n, C_n, D_n , [Journal of Algebra](#) **107**, 466 (1987).

Continued in

k

With ‘large N saturated’ multiplicities and specialisation rules

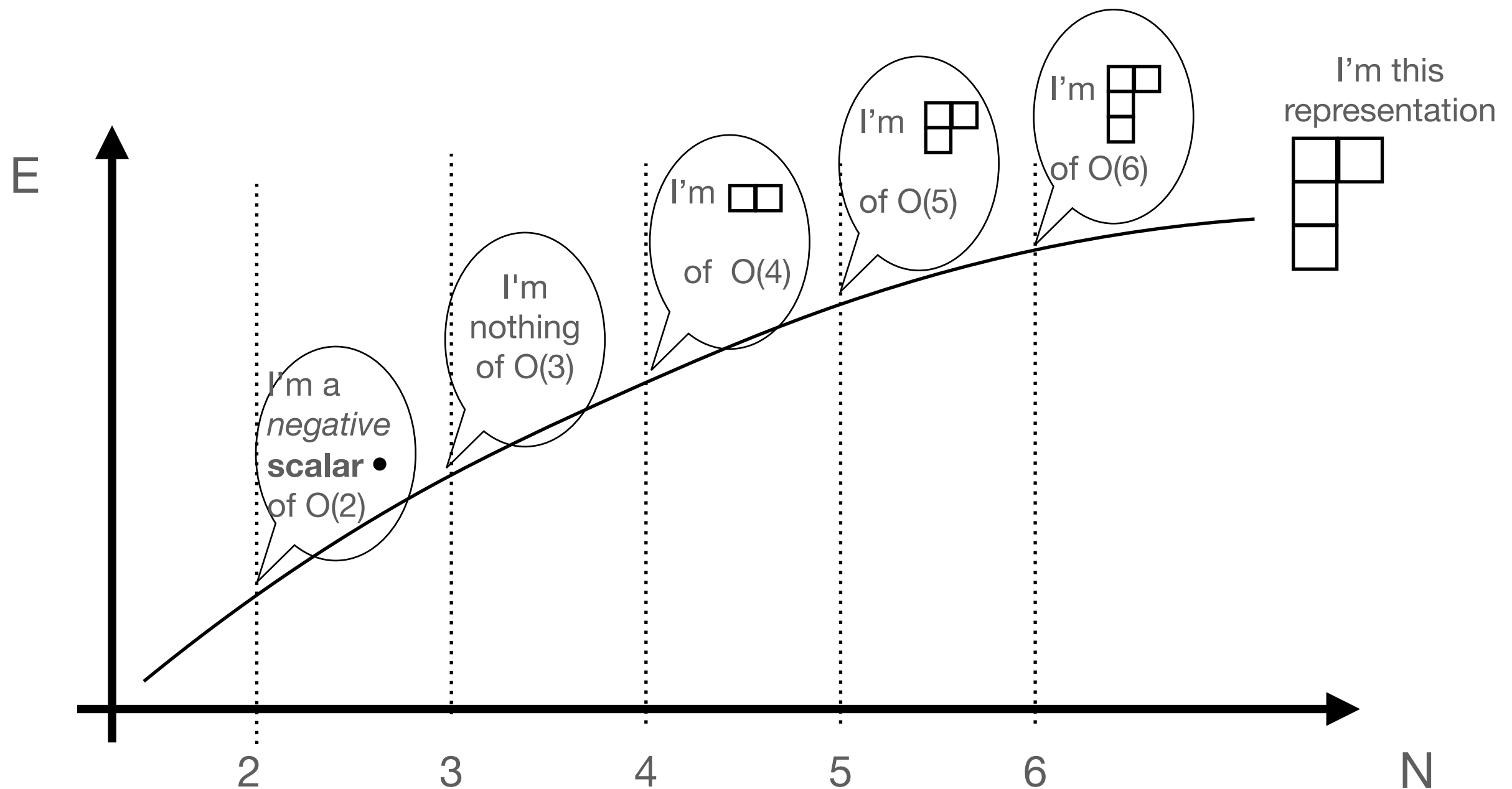
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Example

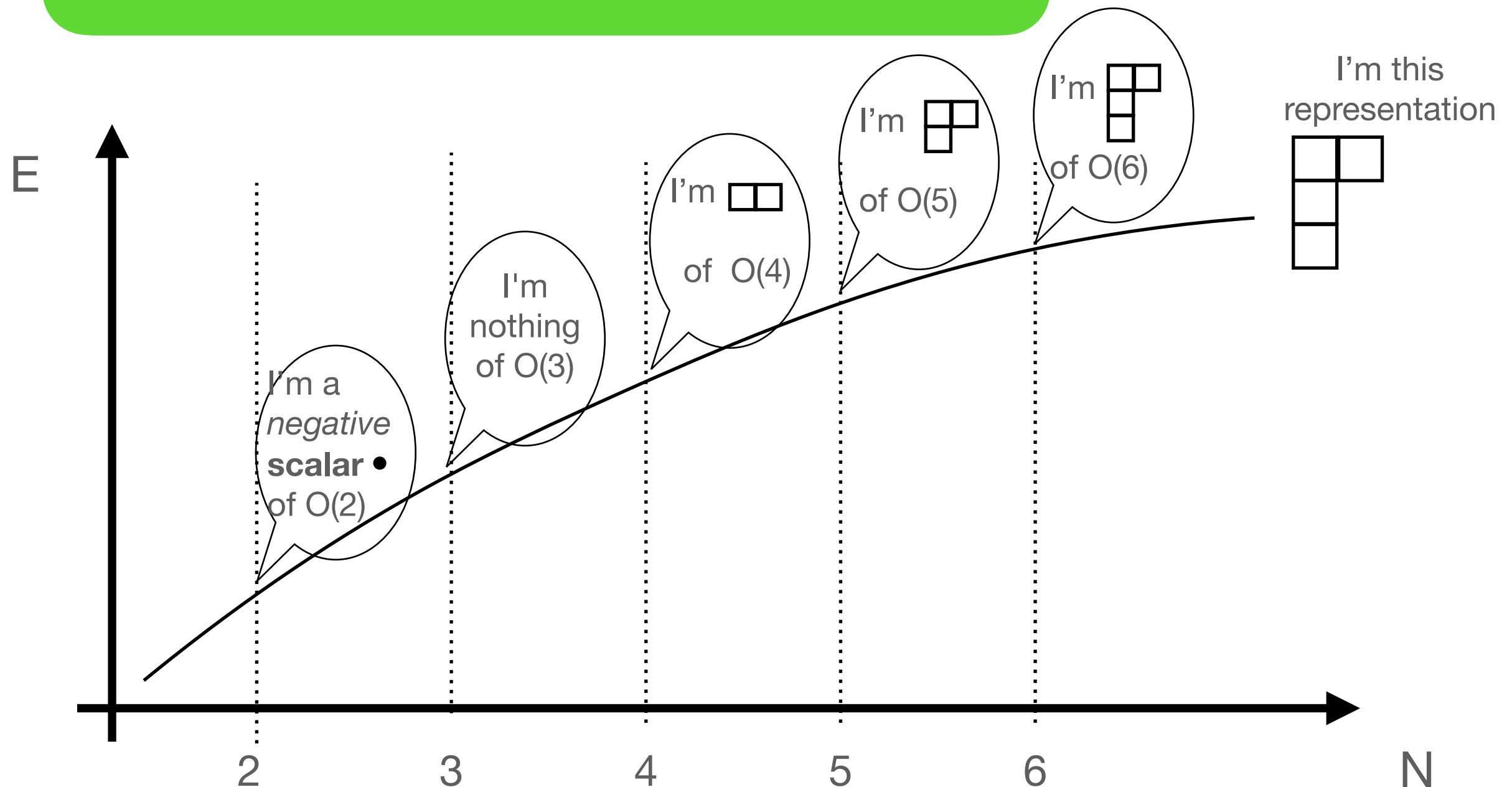
$$Z = \dim \sum_{\Gamma} e^{-E_{\Gamma}(N)/T} + \dots$$



Example

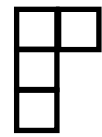
This is what is known in QFT as ‘evanescence’, and we can see there are a number of different types that can happen. (This had not been appreciated before)

+ . . .

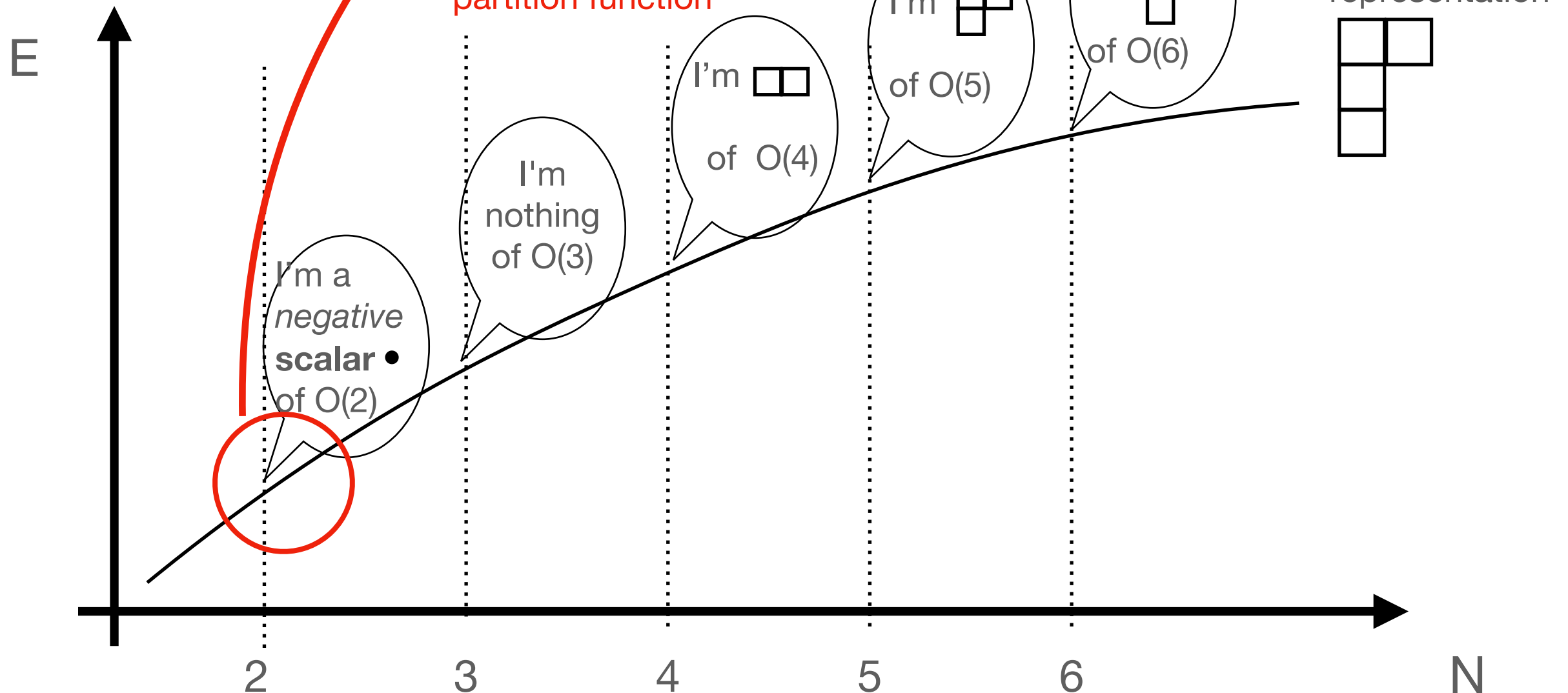


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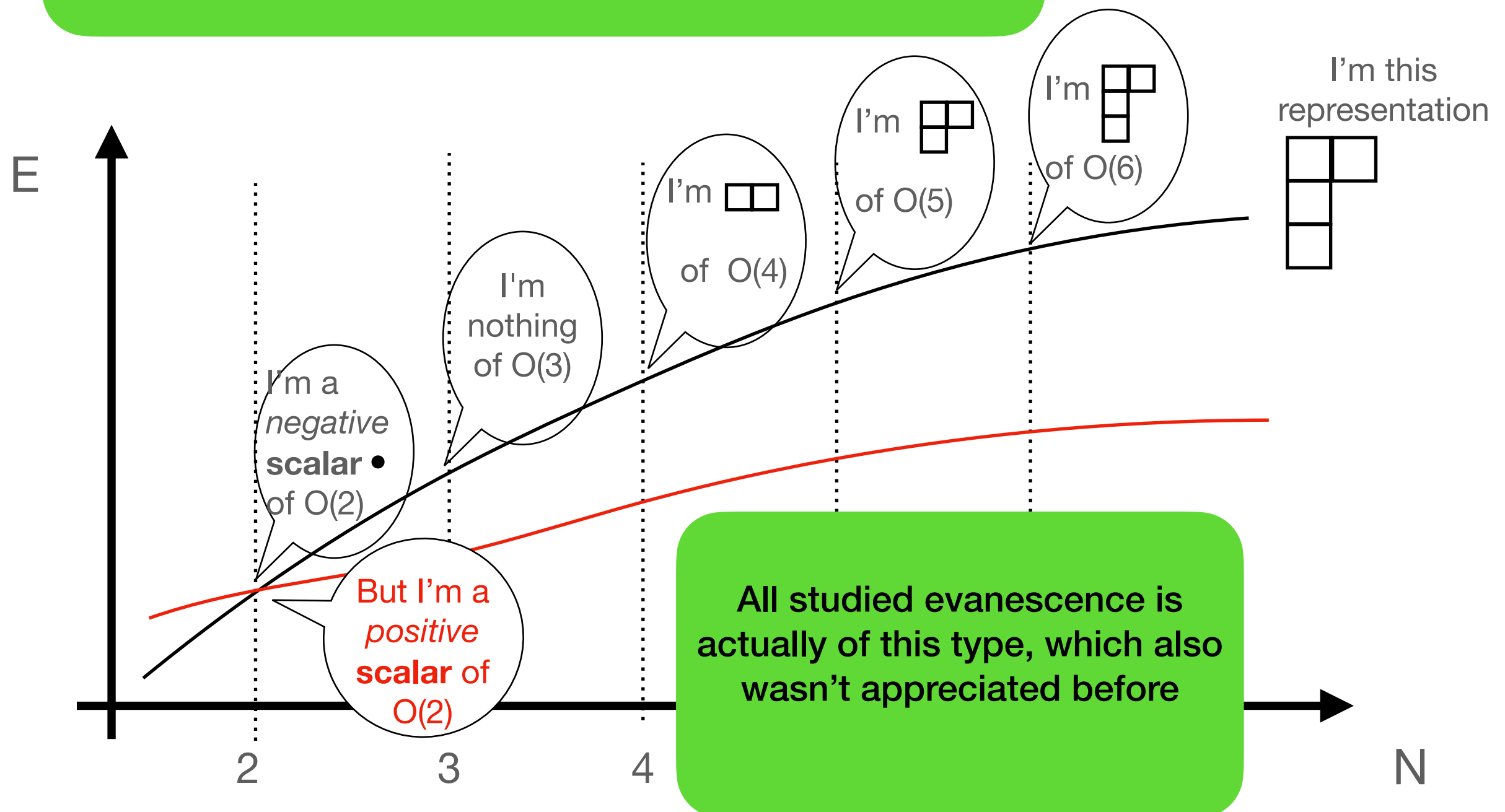
But can't have a negative contribution to the partition function



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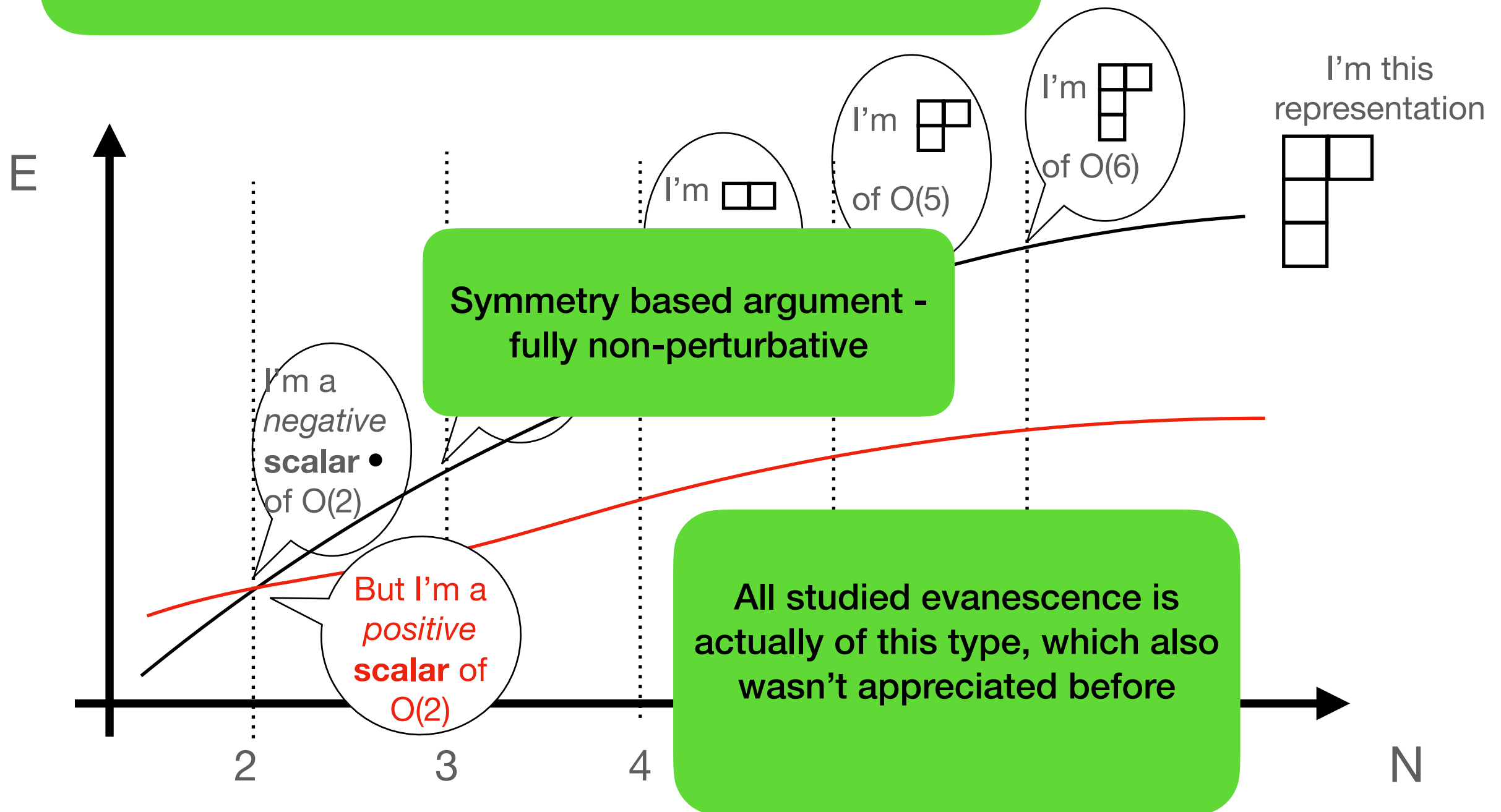
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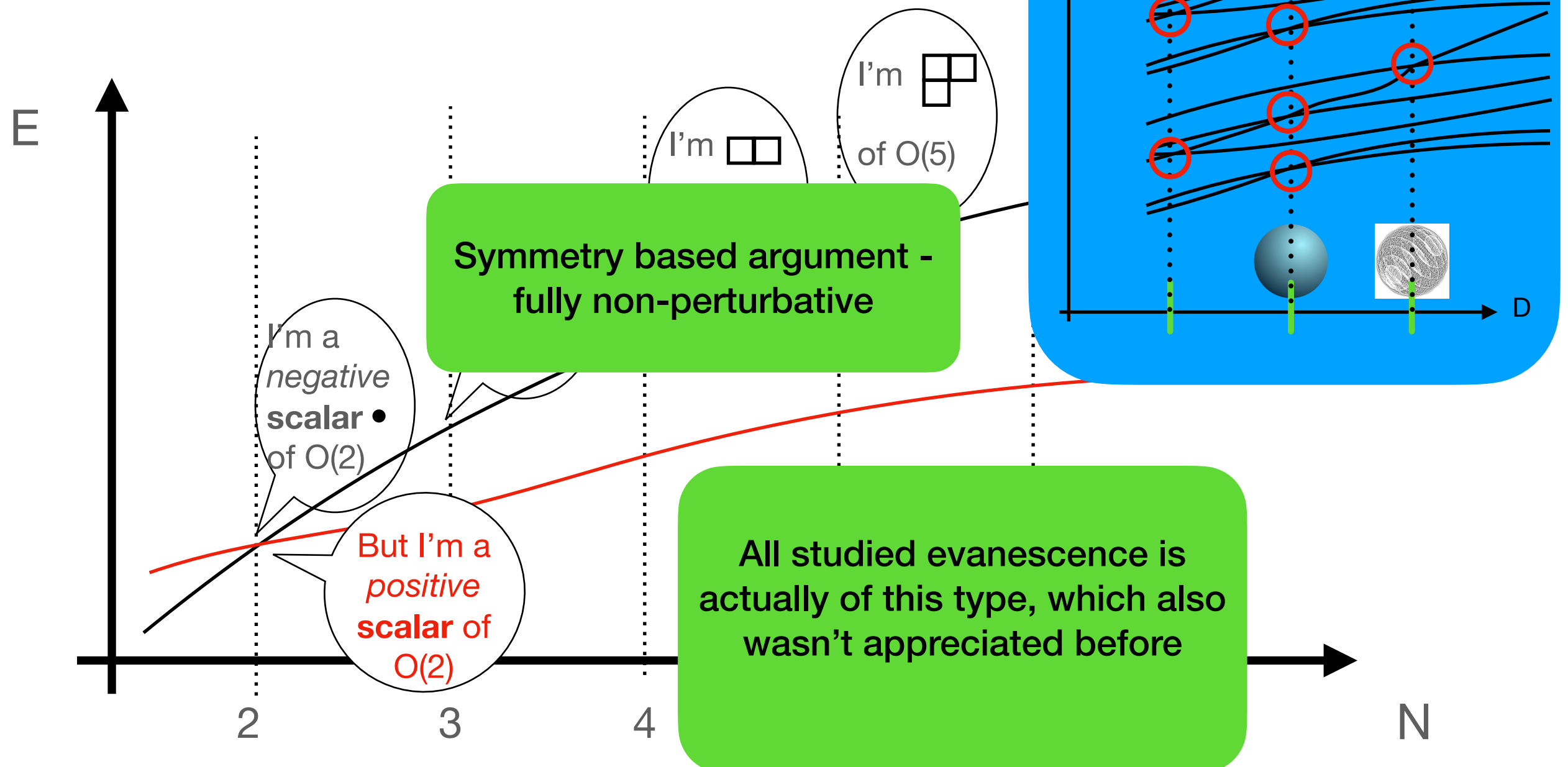
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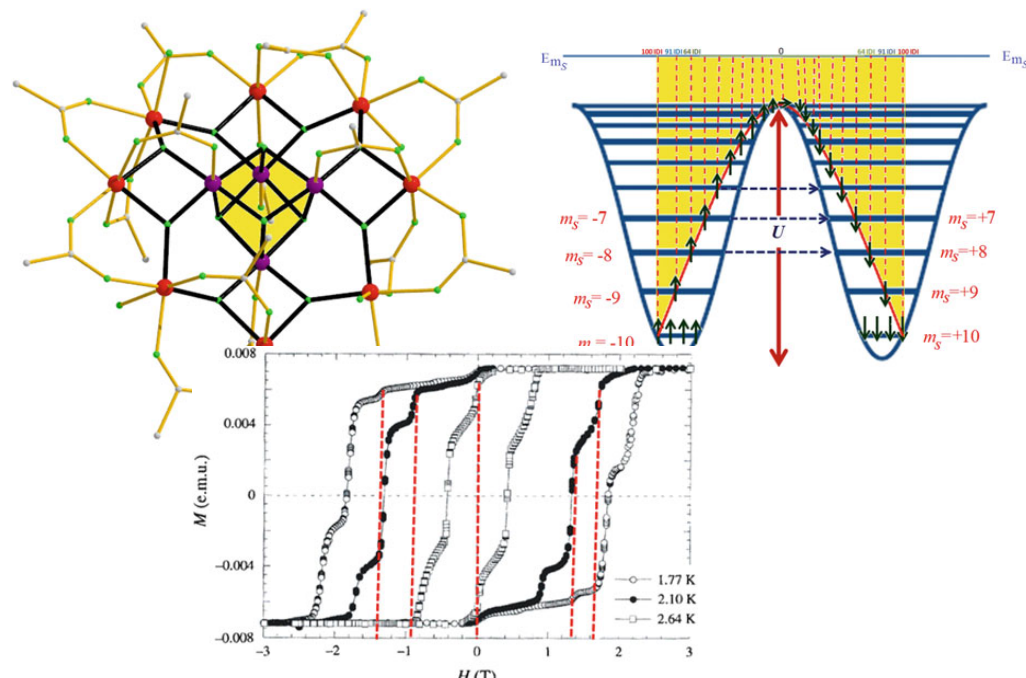


Conclusion

It's quite a surprise that, given how ubiquitous $O(N)$ symmetry is in physics, and how often we play around varying N , there were unnoticed physics consequences of symmetry

Learned something about dim reg

N is an interesting knob to start to investigate mechanisms of unitarity restoration/loss



Unitary
QFT

Non-unitary QFT

